

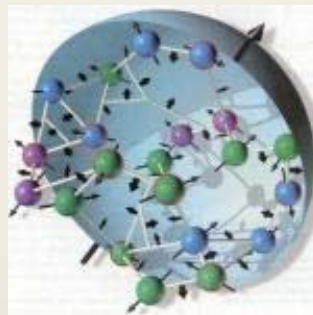


Two Photon Exchange in elastic electron-proton scattering:

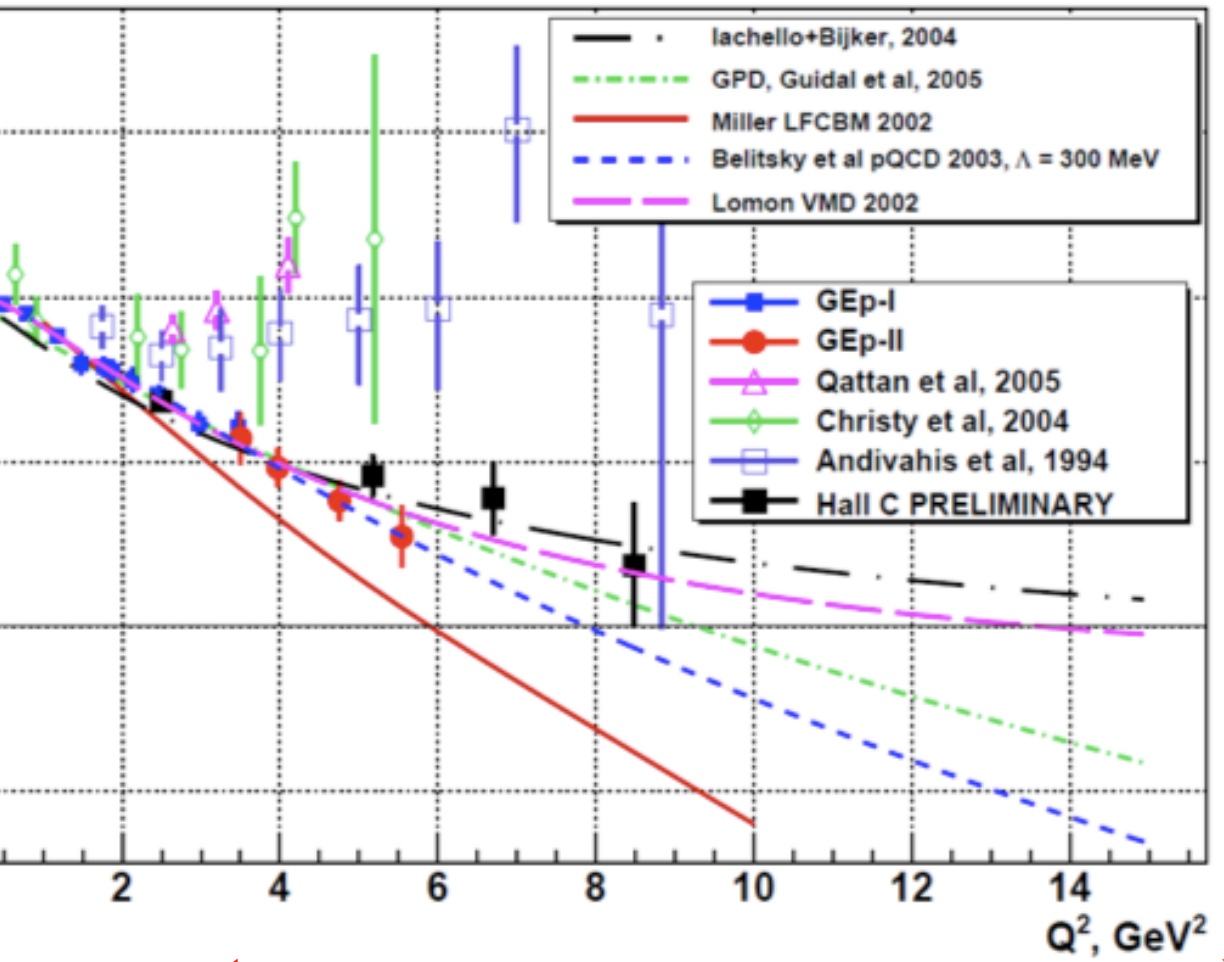
QCD factorization approach Nikolai Kivel

in collaboration with

M. Vanderhaeghen



GEP3 preliminary results: FF ratio



? pQCD in TPE

Theory

Chen, Afanasev et al PRL93(2002)

GPD-model

Blunden Melnitchouk Tjon PRL91(2003)

hadronic model

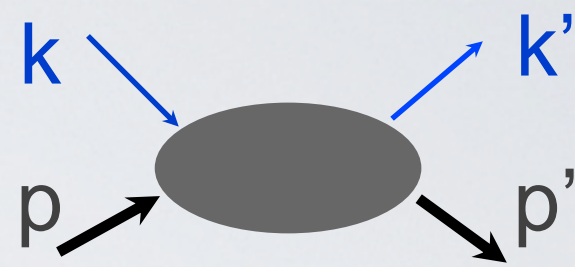
TPE can explain the large discrepancy in the FF extraction

Experiment:

TPE mechanism can be studied experimentally

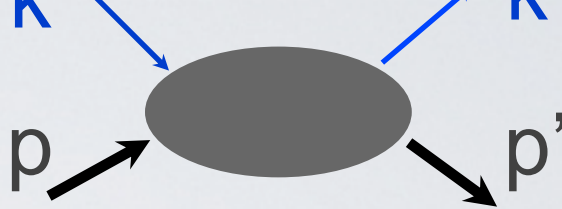
ature

non Vanderhaeghen, 2003



$$A_{ep} = \frac{e^2}{Q^2} \bar{l}' \gamma^\mu l$$
$$\times \bar{N}' \left(\underline{\tilde{G}_M} \gamma^\mu - \underline{\tilde{F}_2} \frac{1}{2M} (p + p')^\mu + \underline{\tilde{F}_3} \frac{1}{4M^2} (p + p')^\mu \gamma \cdot (k + k') \right) N$$

$$-ep = \frac{e^2}{Q^2} \bar{l} \gamma^\mu l$$



$$\bar{N}' \left(\underline{\tilde{G}_M} \gamma^\mu - \underline{\tilde{F}_2} \frac{1}{2M} (p + p')^\mu + \underline{\tilde{F}_3} \frac{1}{4M^2} (p + p')^\mu \gamma \cdot (k + k') \right) N$$

1- and 2 γ - ff's

1 γ exch

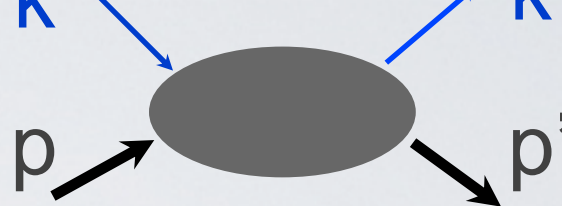
$$\tilde{G}_M = G_M + \delta\tilde{G}_M$$

$$\tilde{F}_2 = F_2 + \delta\tilde{F}_2$$

$$\tilde{F}_3$$

2 γ exch





$$-ep = \frac{e^2}{Q^2} \bar{l} \gamma^\mu l$$

$$\bar{N}' \left(\underline{\tilde{G}_M} \gamma^\mu - \underline{\tilde{F}_2} \frac{1}{2M} (p + p')^\mu + \underline{\tilde{F}_3} \frac{1}{4M^2} (p + p')^\mu \gamma \cdot (k + k') \right) N$$

and 2γ- ff's

1γ exch

$$\tilde{G}_M = G_M + \delta\tilde{G}_M \quad \tilde{F}_2 = F_2 + \delta\tilde{F}_2 \quad \tilde{F}_3$$

large Q² asymptotic (pQCD)

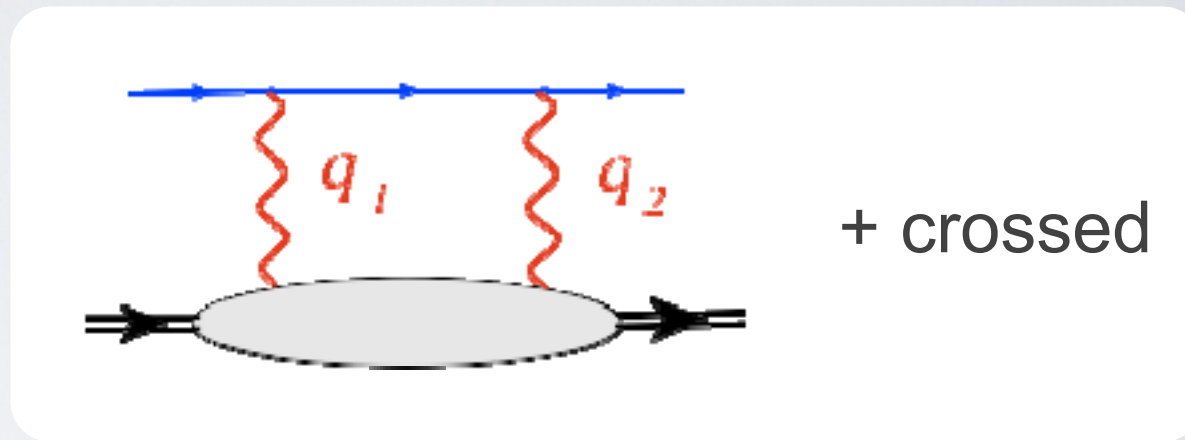
2γ exch

$$G_M \sim F_1 \sim \alpha_S^2 / Q^4$$

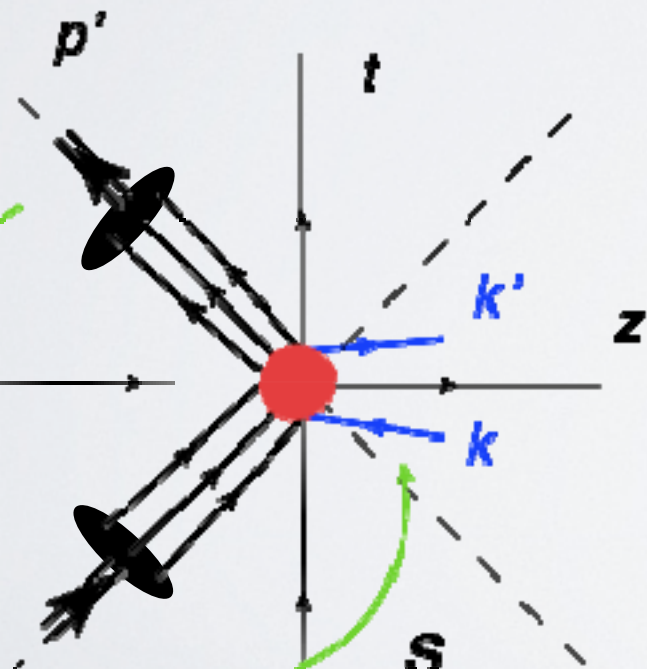
$$F_2 \sim 1/Q^6$$

photons have large
momenta

$$t_2^2 \sim q^2 = (p' - p)^2 \equiv -Q^2$$



Breit frame

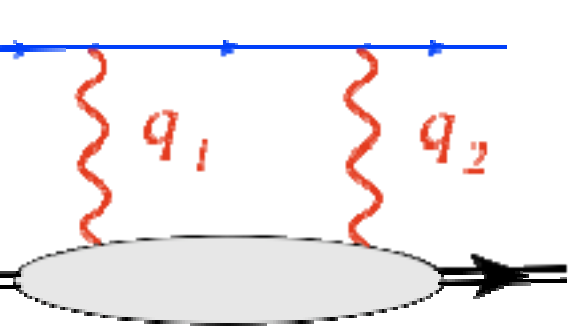


$$\bar{n} = (1, 0, 0, 1) \quad n = (1, 0, 0, -1)$$

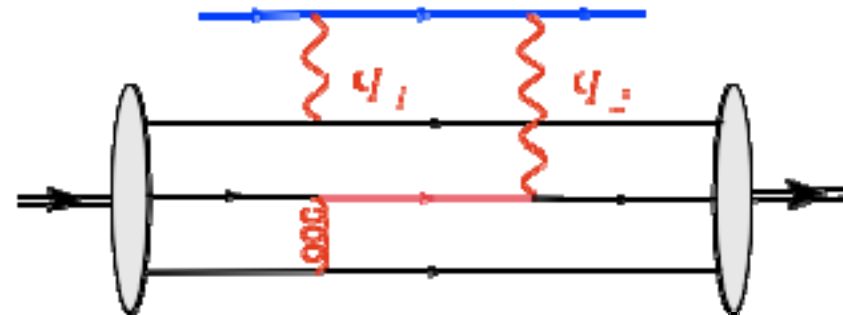
$$p \simeq Q \frac{\bar{n}}{2} \quad p' \simeq Q \frac{n}{2}$$

$$s \rightarrow \infty \quad Q^2/s = z \quad \text{fixed} \quad 0 < z < 1$$

$$k = \frac{1}{2} Q \frac{\bar{n}}{s} + \frac{\bar{z}}{2} Q \frac{n}{s} + k_{\perp} \quad k_{\perp}^2 = -\bar{z}^2 Q^2$$

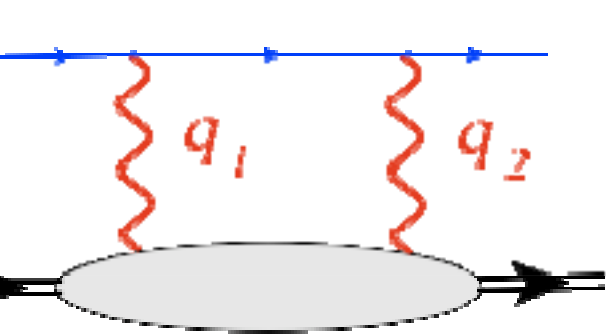


+ crossed =

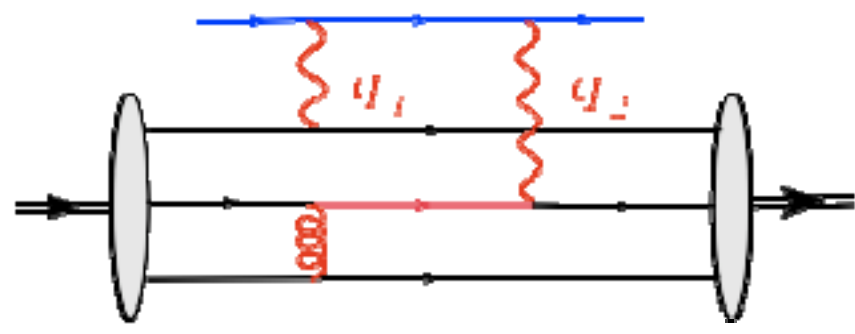


+ others
23 graphs

Lowest order graphs to turn 3 collinear quarks into 3 collinear quarks moving in opposite directions

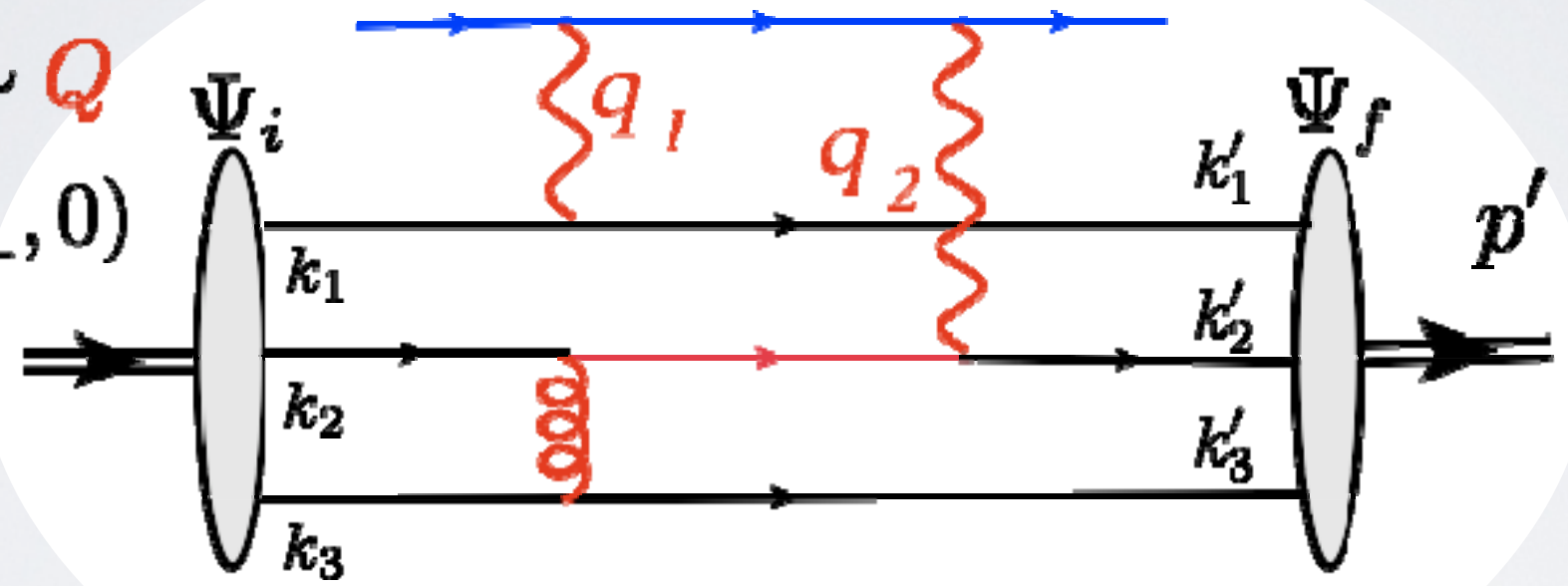


+ crossed =



+ other
23 graphs

$\sim p'_- \sim Q$
 $(p_+, 0_\perp, 0)$



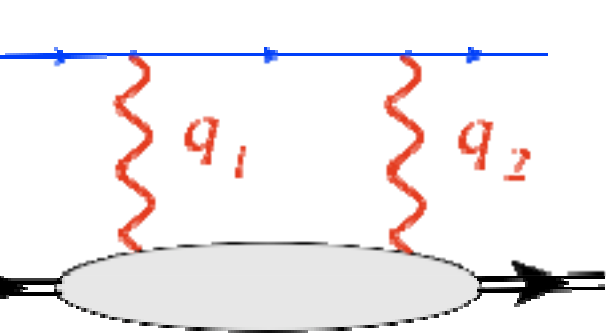
$p' = (0, 0_\perp, p'_-)$

$= x_i p + k_{i\perp}$

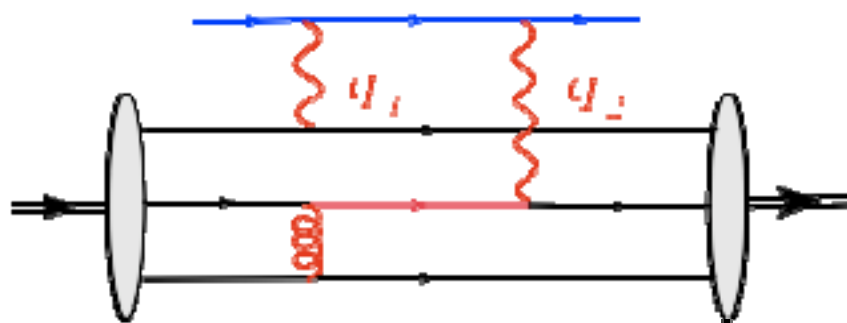
$k_{i\perp} \sim k'_{i\perp} \sim \Lambda_{QCD} \ll Q$

$= x'_i p + k'_{i\perp}$

\int



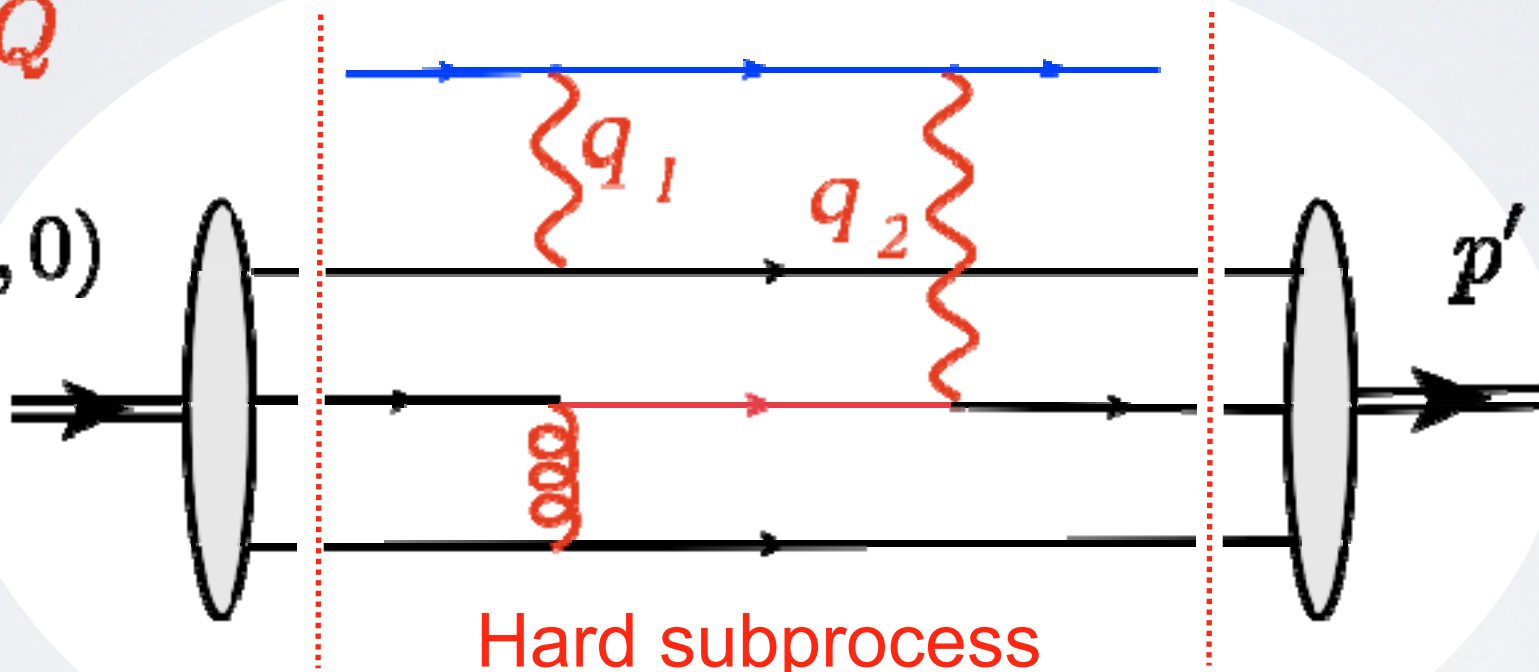
+ crossed =



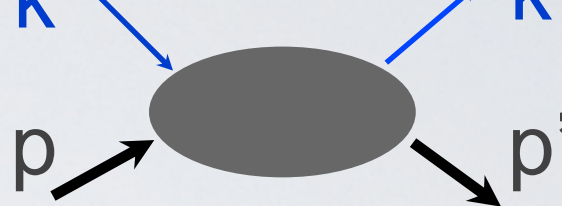
+ other
23 graphs

$$p'_- \sim Q$$

$$(p_+, 0_\perp, 0)$$



$$\Phi(x_i) * T(x_i, x'_i | z, Q^2) * \Phi(x'_i)$$



$$-ep = \frac{e^2}{Q^2} \bar{l} \gamma^\mu l$$

$$\bar{N}' \left(\underline{\tilde{G}_M} \gamma^\mu - \underline{\tilde{F}_2} \frac{1}{2M} (p + p')^\mu + \underline{\tilde{F}_3} \frac{1}{4M^2} (p + p')^\mu \gamma \cdot (k + k') \right) N$$

and 2γ- ff's

1γ exch

$$\tilde{G}_M = G_M + \delta\tilde{G}_M \quad \tilde{F}_2 = F_2 + \delta\tilde{F}_2 \quad \tilde{F}_3$$

large Q² asymptotic (pQCD)

2γ exch

$$G_M \sim F_1 \sim \alpha_S^2 / Q^4$$

$$\delta\tilde{G}_M \sim \frac{(PK)}{M^2} \tilde{F}_3 \sim \alpha_{em} \alpha_S / Q^4$$

CD FF's:

$$** \equiv \int d^3x d^3x' \delta(1 - \sum x_i) \delta(1 - \sum x'_i)$$

$$\left. \begin{matrix} M \\ \tilde{F}_3 \end{matrix} \right\} = \mathcal{K} \left\{ \begin{matrix} T_M \\ T_3 \end{matrix} \right\} (x_i, x'_i | z) ** \{ Q_u Q_d [\varphi_N(x'_i) \varphi_N(x_i) + \dots] + Q_u^2 [\dots] + Q_u Q_d [\dots] \}$$

$$\mathcal{K} = - \frac{\alpha_{em} \alpha_s(\mu^2)}{Q^4} \left(\frac{2\pi}{3} \right)^2 \frac{Q^2 (s - u)}{s^2}$$

$$\left. \right\} (x_i, x'_i) = \left\{ \begin{matrix} 4x_2 x'_2 \\ 2(x_2 + x'_2 - 2x_2 x'_2) \end{matrix} \right\} \left[\frac{1}{x'_1 x'_2 (1 - x'_2)} \frac{1}{x'_2 - x_2 + x_2 (1 - x'_2) Q^2/s + i0} \times (x \leftrightarrow x') \right]$$

pole in the lepton propaga
 nontrivial Im part

All integrals are IR-finite

neon DA

$$\varphi_N(132) \simeq 120 f_N x_1 x_2 x_3 (1 + \tau_- (x_1 - x_2) + \tau_+ (1 - 3x_3) + \dots)$$

- assumption: drop the highest conformal moments

uncertainties: renorm. scale (NLO) $\alpha_S(\mu)$

helicity flip FF: $\delta \tilde{G}_E = \delta \tilde{F}_1 - \tau \delta \tilde{F}_2$

3 non-perturbative parameters

models for DA :

	f_N (10^{-3} GeV ²)	τ_-	τ_+
CD SR (1988), COZ	5.0 ± 0.5	4.0 ± 1.5	1.1 ± 0.5
LCSR (2006)			

$$= \underline{G_M^2 + \frac{\epsilon}{\tau} G_E^2} + 2G_M \mathcal{R} \left(\delta\tilde{G}_M + \epsilon \frac{\nu}{M^2} \tilde{F}_3 \right) + 2\frac{\epsilon}{\tau} G_E \mathcal{R} \left(\delta\tilde{G}_E + \frac{\nu}{M^2} \tilde{F}_3 \right)$$

Born app

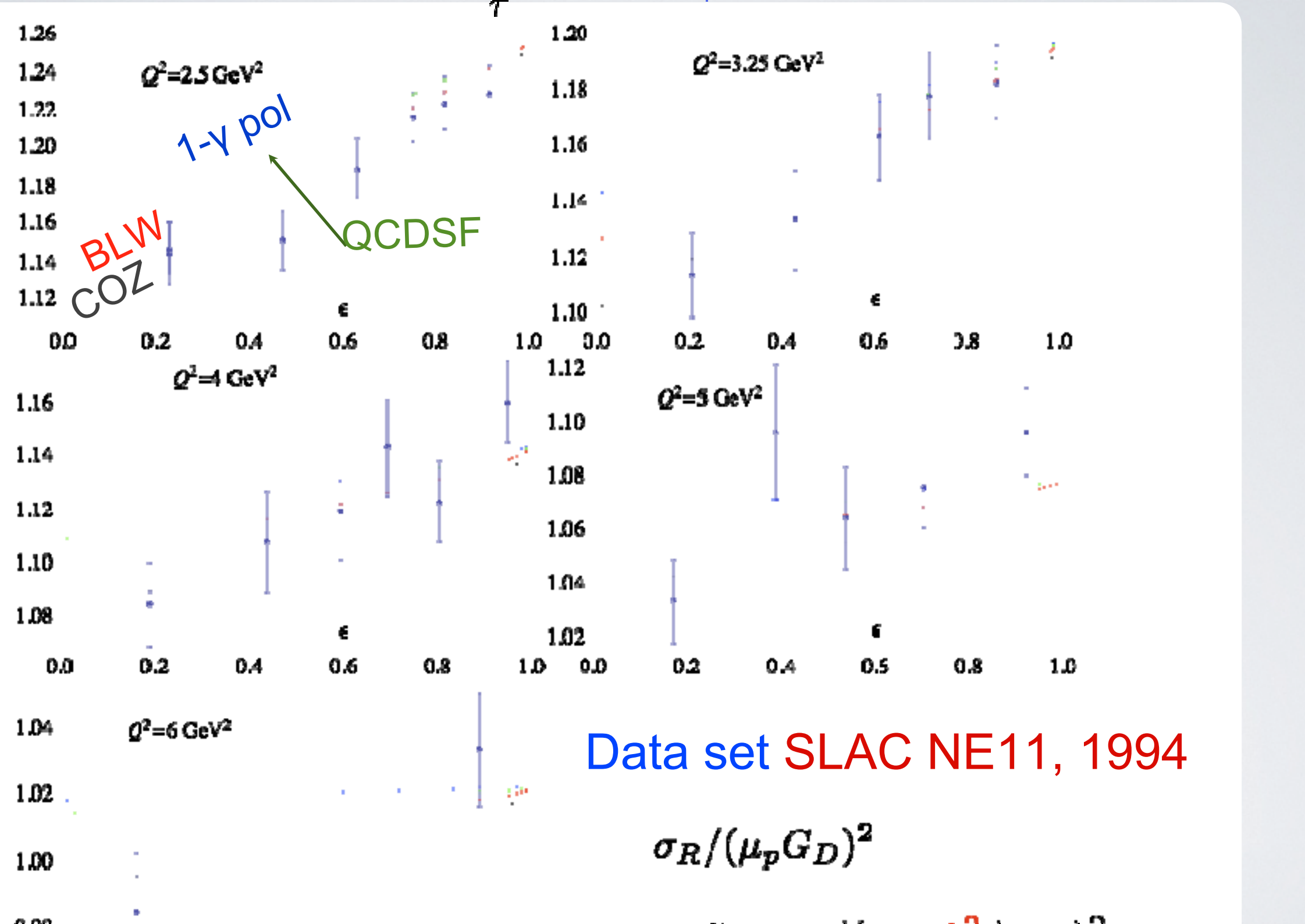
Re part of TPE

Empirical fit for the ratio and form factor:

$$(\mu_p G_E / G_M)_{exp} = 1 - 0.13(Q^2 - 0.04) \quad 0.5\text{GeV}^2 < Q^2 < 5.6\text{GeV}^2$$

JLab Hall A, 2002

$$G_M = \frac{\mu_p \theta(0.1 < Q^2 < 6)}{(1 + 0.116Q + 2.874Q^2 + 0.241Q^3 + 1.00Q^4 + 0.345Q^5)}$$



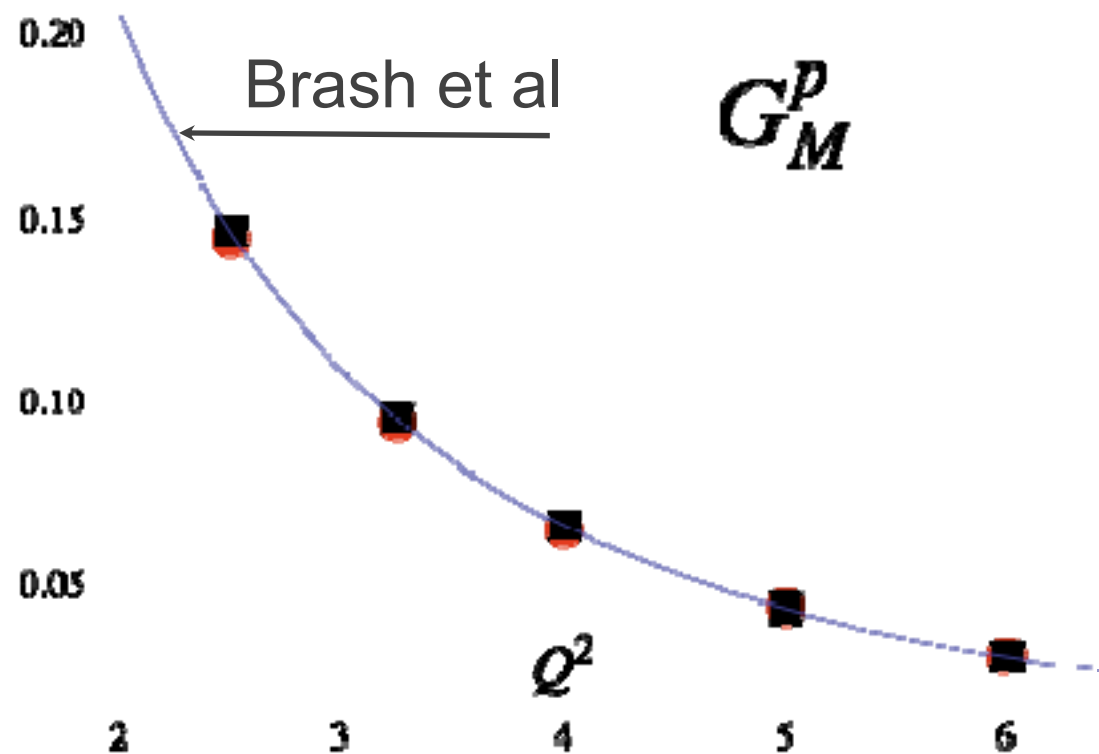
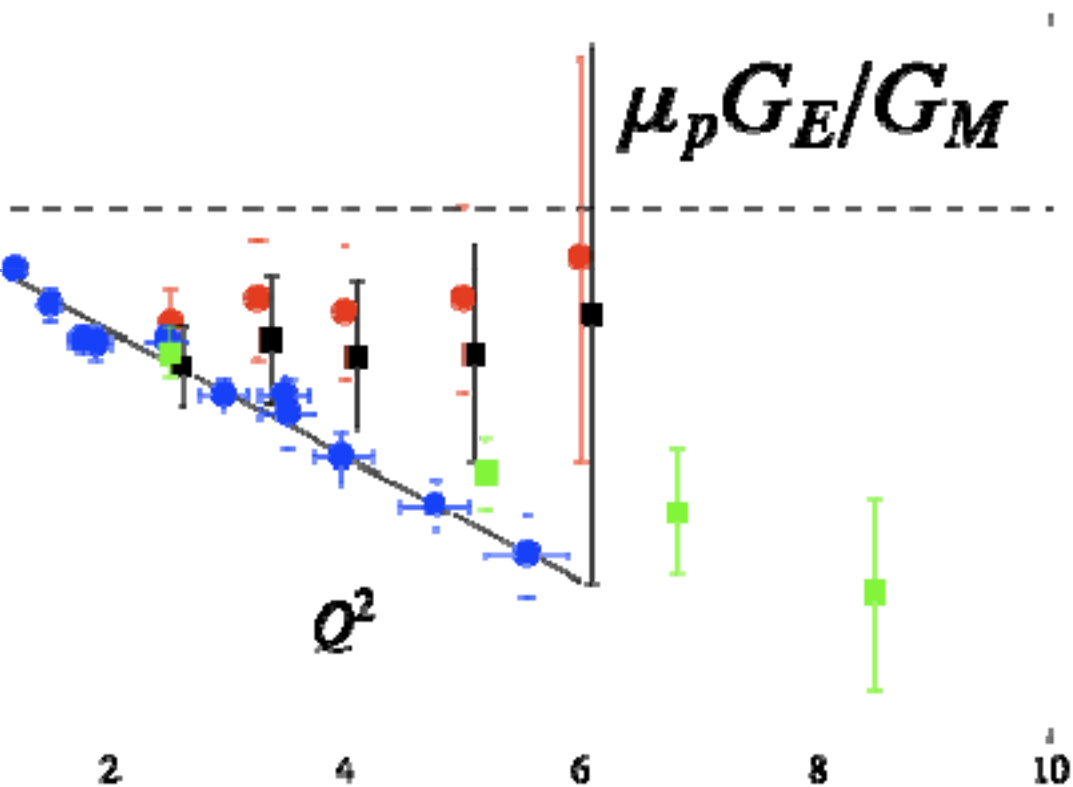
cross section $\sigma_R = G_M^2 + \frac{\epsilon}{\tau} G_E^2 + \Delta\sigma_{2\gamma}$

Data set SLAC NE11, 1994

● BLW

■ COZ

NK, Vanderhaeghen, in prep.



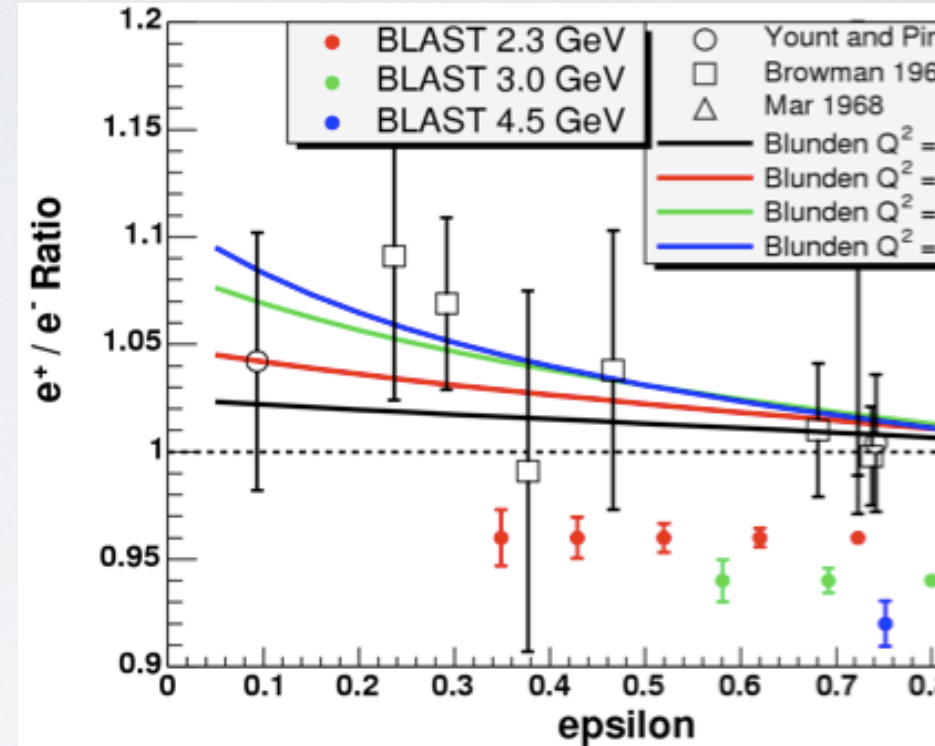
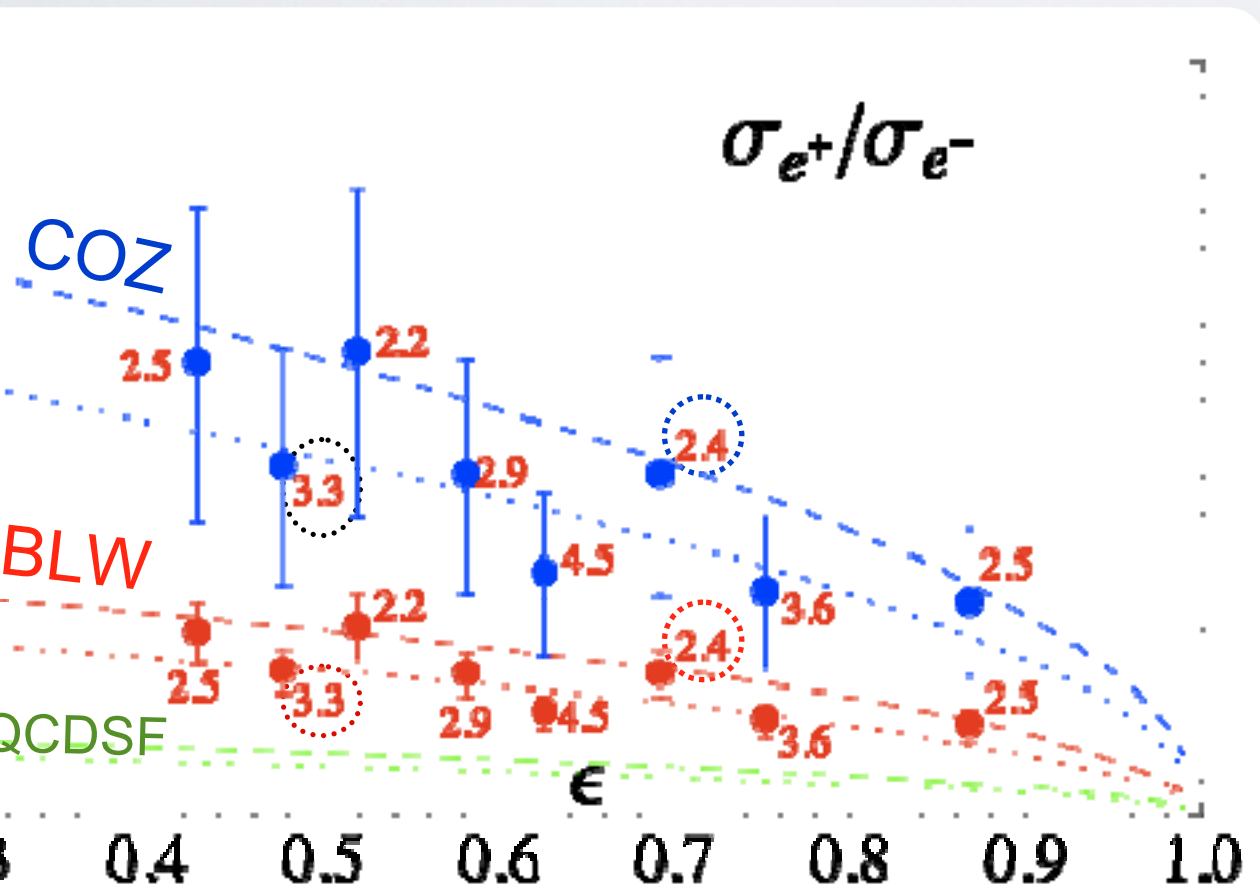
Future experiments:

☀ JLab Class PR-07-005 $Q^2=.5-2.5$ GeV² $\epsilon=.1$

☀ VEPP-3 Novosibirsk $Q^2=1.6$ GeV² $\epsilon=.44, .9$

☀ Olympus@Desy $Q^2=.8-4.5$ GeV² $\epsilon=.4-.9$

CD estimate

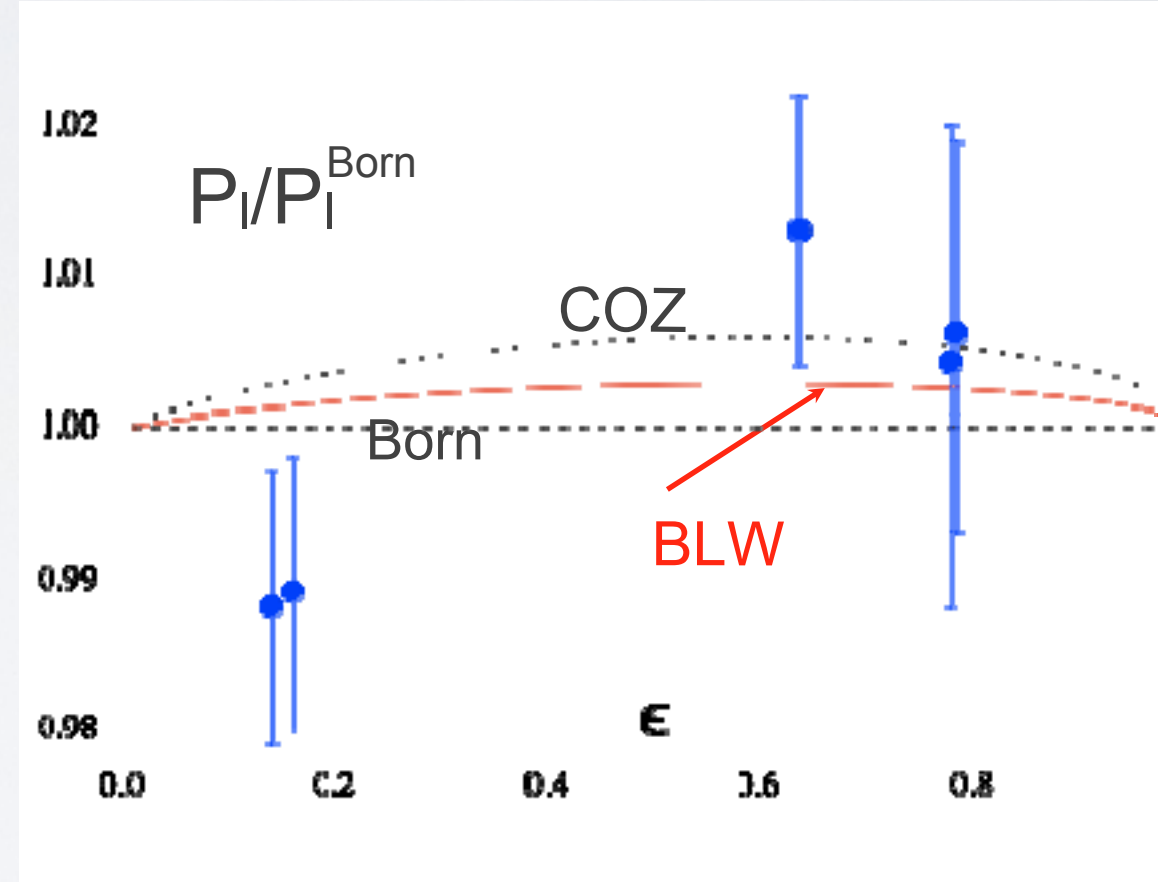
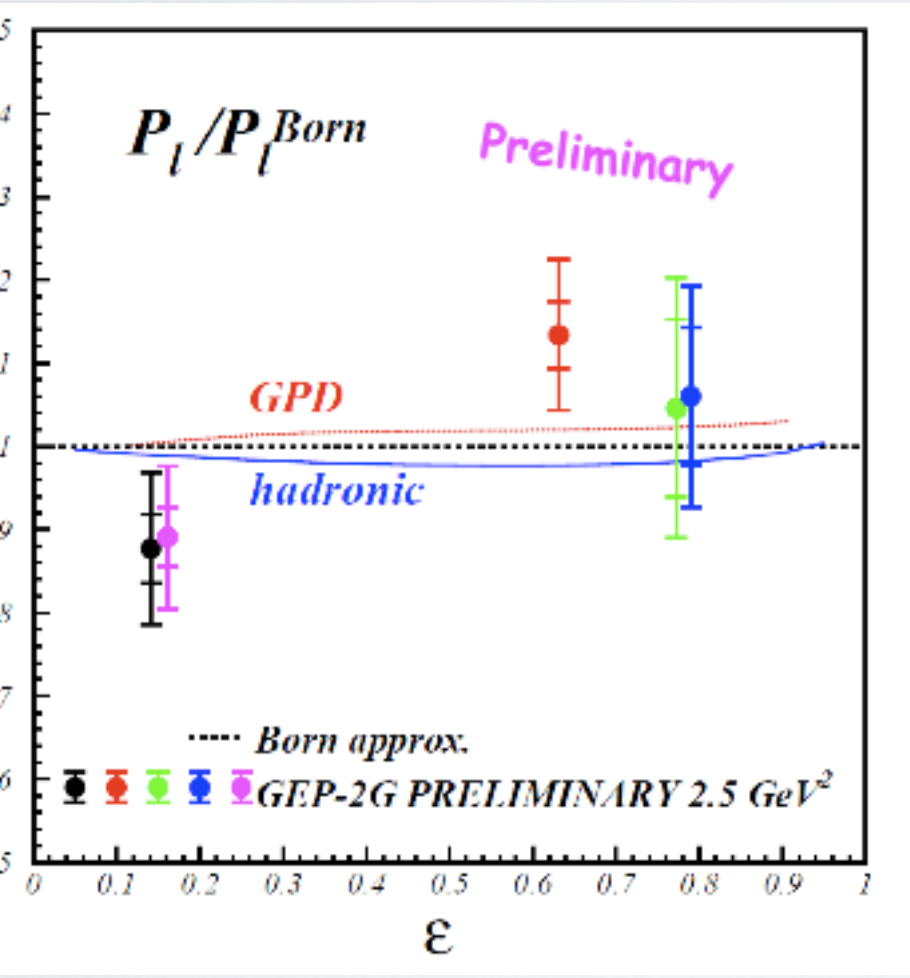


Expected accuracy <

$$P_l = \sqrt{1 - \epsilon^2} \frac{1}{\sigma_H} \left\{ G_M^2 + 2 G_M \mathcal{R} \left(\delta \ddot{G}_M + \frac{\epsilon}{1 + \epsilon} \frac{\nu}{M^2} \ddot{F}_3 \right) + \mathcal{O}(\epsilon^4) \right\}$$

Hall C E-04-119, preliminary
 5 GeV² $\epsilon = .14, .63, .785$

pQCD, different models



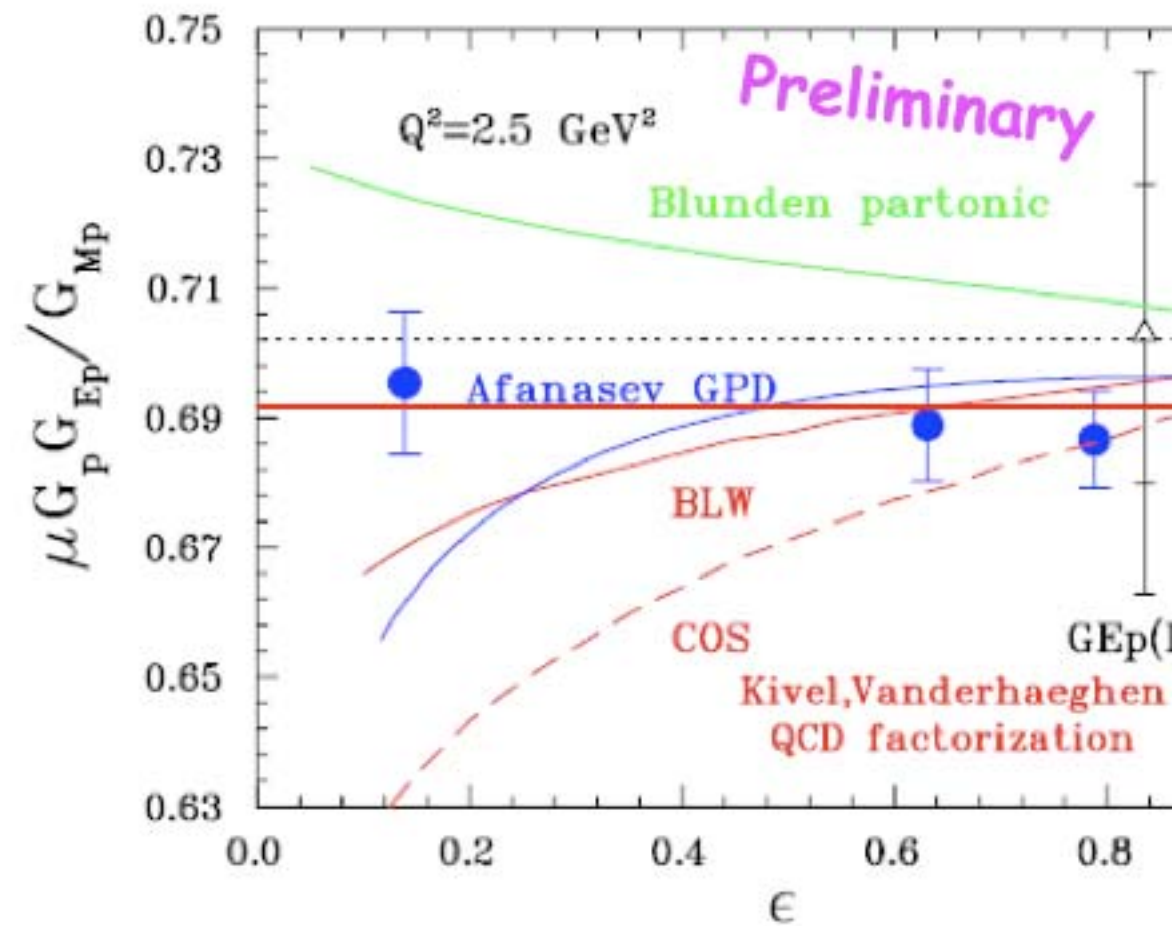
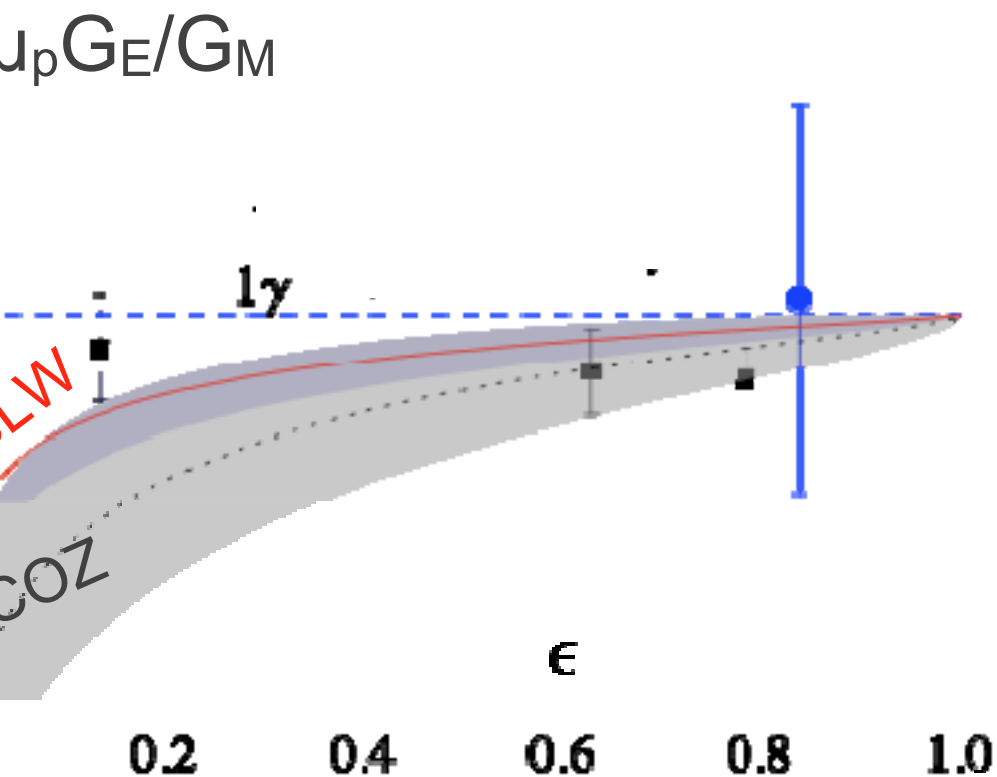
$$= -\sqrt{\frac{2\varepsilon(1-\varepsilon)}{\tau}} \frac{1}{\sigma_R} \left\{ G_E G_M + G_E \mathcal{R}(\delta \tilde{G}_M) + G_M \mathcal{R}(\delta \tilde{G}_E + \frac{\nu}{M^2} \tilde{F}_3) + \mathcal{O}(\varepsilon^4) \right\}$$

$$G_E \sim \lambda \delta \tilde{G}_M \quad \lambda = 0, 1$$

QCD, different models

JLab Hall C E-04-119, preliminary

$Q^2=2.5 \text{ GeV}^2$ $\varepsilon=.14,.63,.785$



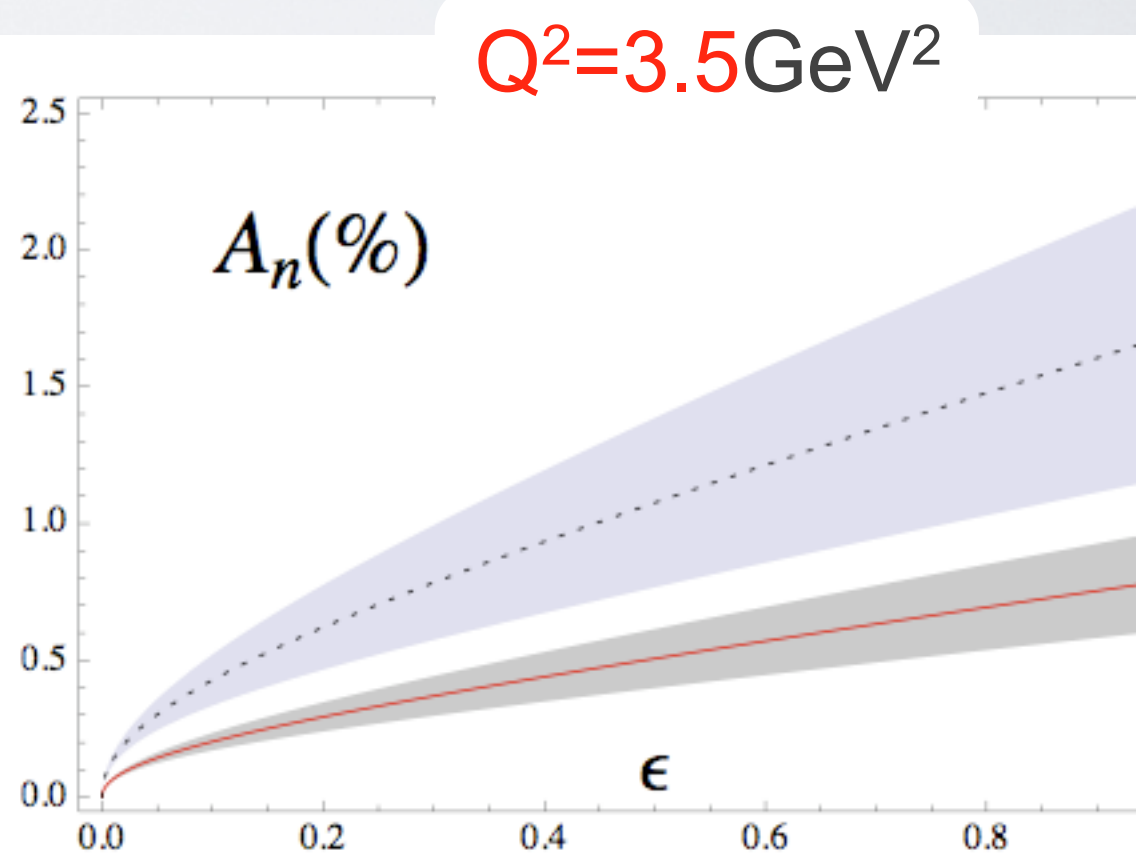
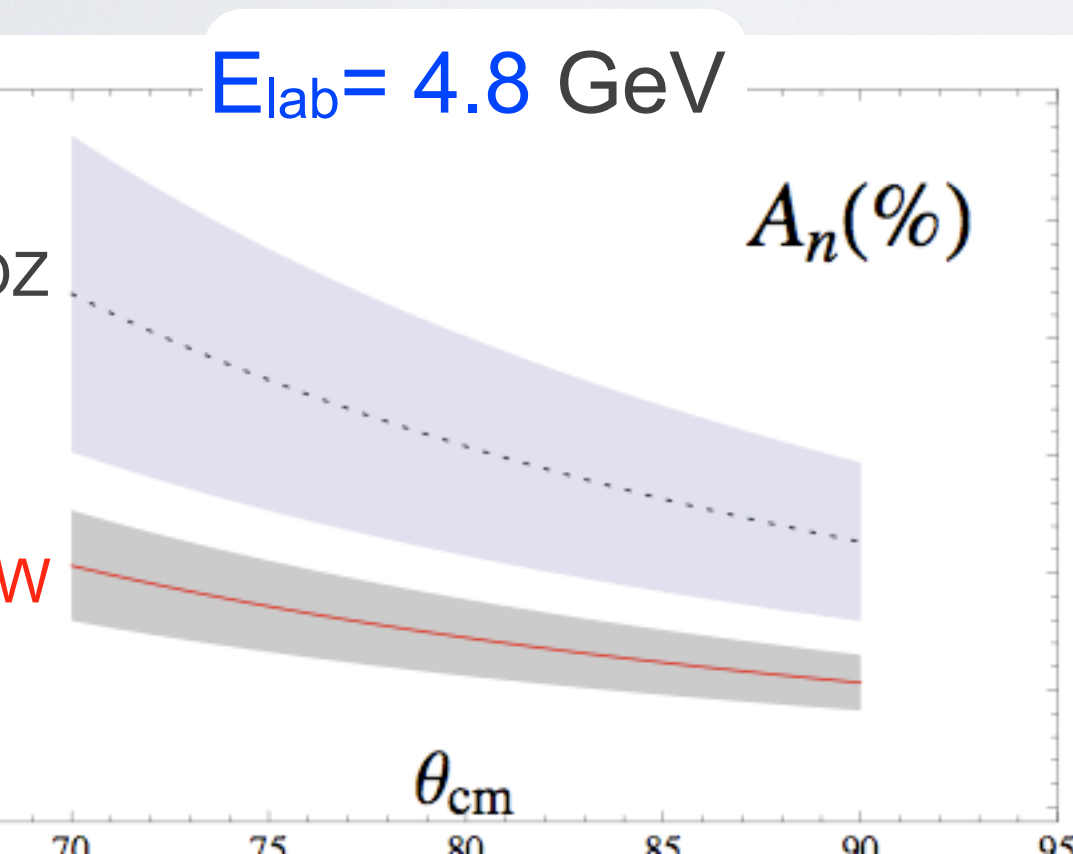
Target Normal Spin Asymmetry

$$= \sqrt{\frac{2\varepsilon(1+\varepsilon)}{\tau}} \frac{1}{\sigma_R} \left(-G_M \text{Im} \left\{ \delta\tilde{G}_E + \frac{(K \cdot P)}{M^2} \tilde{F}_3 \right\} + G_E \text{Im} \left\{ \delta\tilde{G}_M + \frac{2\varepsilon}{1+\varepsilon} \frac{(K \cdot P)}{M^2} \tilde{F}_3 \right\} \right)$$

$\sigma_L/\sigma_T \simeq 0.35$ E-02-013 Hall-C, preliminary

$G_M \simeq \mu_n G_D$ CLASS, 2009 $Q^2=1.-4.5 \text{ GeV}^2$

Hall-A, E-05-015 $E_{\text{lab}}=3.3, 5.3 \text{ GeV}$



hard QCD contribution in TPE at large Q^2 is not small and must be considered accurately

TPE at large Q^2 depends on the structure of target and can provide us information about nucleon DA

experiments at higher Q^2 where the theory works better is a good opportunity to check pQCD predictions

large power corrections and/or NLO can change optimistic picture. This question must be studied.